

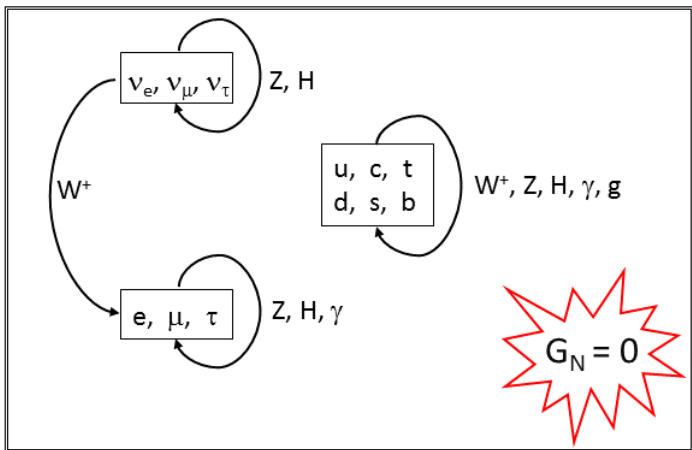
Symmergent Gravity and Ultraviolet Stability

Durmuş Demir

İZMİR INSTITUTE OF TECHNOLOGY

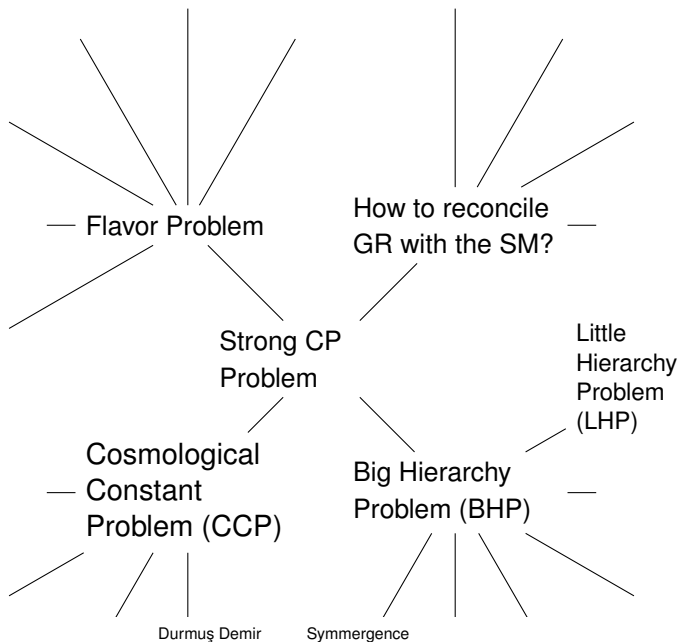
HEPAC (February 08, 2018), İstanbul, Turkey

Working Model

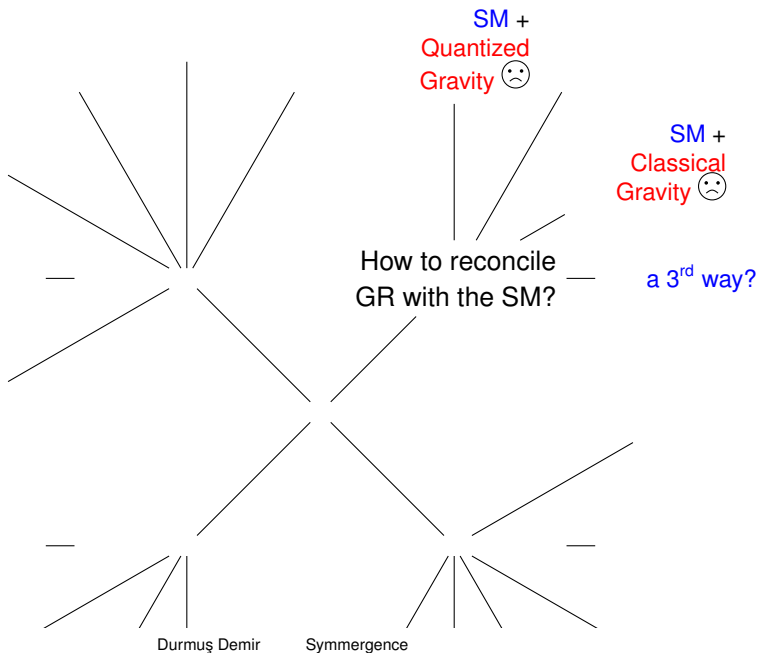


Standard Model = a renormalizable QFT of the
EM (γ) + Weak (W, Z, H) + Strong (g)

Gordian Knots In The SM



Gravity Knot



3rd Way

3rd Way = SM + Emergent Gravity

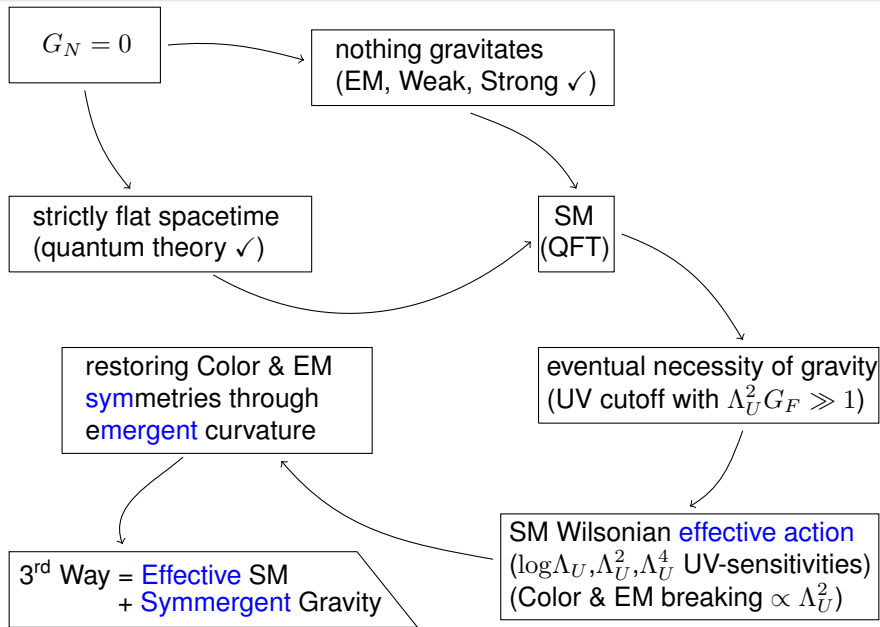
Sakharov's Induced Gravity

Verlinde's Entropic Gravity

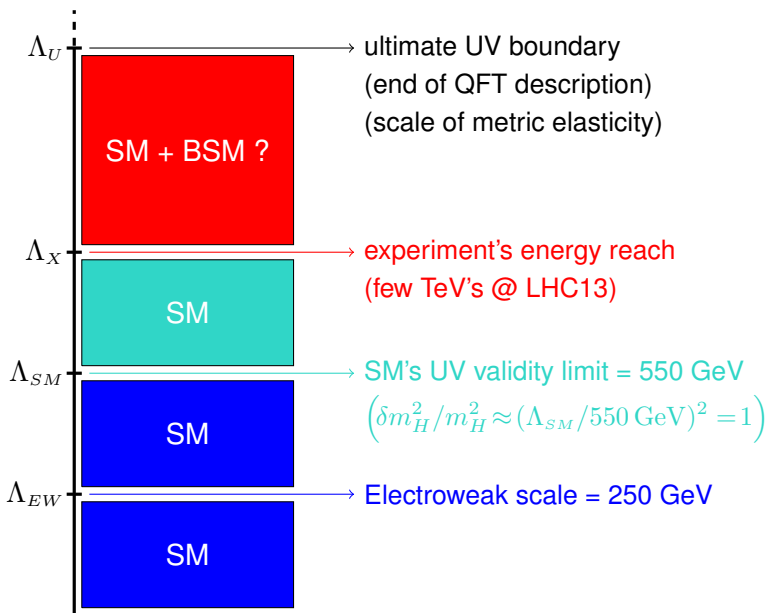
Raamsdonk's Entanglement Gravity

► Some of these can be in tension with Weinberg-Witten theorem !

3rd Way – A New Kind



Scales & QFTs



A Prudent QFT Setup

- ▶ Assume that the **SM**, which has already crossed Λ_{SM} , **holds good** all the way up to Λ_U .
- ▶ Determine then possible **BSM** physics by physical **consistency**.

SM Effective Action

The diagram illustrates the energy scales and fields in the SM Effective Action. A horizontal line represents the energy axis, with two vertical tick marks labeled Λ_X and Λ_U . Above the line, the energy range $(E < \Lambda_X)$ is associated with the Standard Model fields $\psi_{SM}(E)$ and the set $\{\text{leptons, quarks, } H, V_\mu\}$. The energy range $(\Lambda_X < E < \Lambda_U)$ is associated with the fast fields $\psi_{SM}^{(fast)}(E)$. Below the line, the effective action is given by the integral:

$$e^{iS_{eff}(\eta, \psi_{SM})} = \int_{[\Lambda_X]}^{[\Lambda_U]} \mathcal{D}\psi_{SM}^{(fast)} e^{iS(\eta, \psi_{SM}, \psi_{SM}^{(fast)})}$$

Arrows indicate that the fields $\psi_{SM}(E)$ and the set $\{\text{leptons, quarks, } H, V_\mu\}$ are associated with the energy range $(E < \Lambda_X)$, and the fast fields $\psi_{SM}^{(fast)}(E)$ are associated with the energy range $(\Lambda_X < E < \Lambda_U)$. The flat metric $\eta_{\mu\nu}$ is associated with the action $S(\eta, \psi_{SM}, \psi_{SM}^{(fast)})$.

SM Effective Action

$$S_{eff} = S_{tree}(\eta, \psi_{SM})$$

$$+ \underbrace{\delta S_{log}\left(\eta, \psi_{SM}, \log \frac{\Lambda_X}{\Lambda_U}\right)}_{\text{at the SM scale}}$$

$$+ \underbrace{\delta S_{power}\left(\eta, H, V_\mu, \Lambda_U^2 - \Lambda_X^2, \Lambda_U^2 + \Lambda_X^2\right)}_{\text{at the UV scale}}$$

$$\Lambda_U^2 + \Lambda_X^2$$

$$\frac{\Lambda_X}{\Lambda_U} = \text{hierarchy}$$

$$\Lambda_U^2 - \Lambda_X^2$$

UV-Born Parts

$$\delta S_{power} = \delta S_O + \delta S_H + \delta S_V$$

$$\int d^4x \sqrt{-\eta} c_V \left(\frac{\Lambda_X}{\Lambda_U} \right) (\Lambda_U^2 - \Lambda_X^2) \text{Tr}[V_\mu V^\mu]$$

$$-\int d^4x \sqrt{-\eta} c_H \left(\frac{\Lambda_X}{\Lambda_U} \right) (\Lambda_U^2 - \Lambda_X^2) H^\dagger H$$

$$-\int d^4x \sqrt{-\eta} c_O \left(\frac{\Lambda_X}{\Lambda_U} \right) (\Lambda_U^4 - \Lambda_X^4)$$

UV-Caused Knots

$$\delta S_{\text{power}} = \delta S_O + \delta S_H + \delta S_V$$

no problem
in flat ST

**big
hierarchy
problem**



hard UV masses to all
 $U(1)_Y, SU(2)_L, SU(3)_C$
gauge bosons

**explicit breaking of
Color & EM**



► How to untangle these knots?

How To Incorporate Gravity Into SM Effective Action?

$$S_{eff}(\eta, \psi_{SM}, \Lambda_X, \Lambda_U)$$

$$\eta_{\mu\nu} \iff g_{\mu\nu}$$

$$S_{eff}(g, \psi_{SM}, \Lambda_X, \Lambda_U)$$

“curvature” can be **added** by hand

$$S_{eff}(g, \psi_{SM}, \Lambda_X, \Lambda_U) - \int d^4x \sqrt{-g} \left(\tilde{M}^2 R(g) + \tilde{V} + \dots \right)$$

incalculable, arbitrary constants 😞
(loops have already been used up)

How To Incorporate Gravity Into SM Effective Action?

$$S_{eff}(\eta, \psi_{SM}, \Lambda_X, \Lambda_U)$$

$$\eta_{\mu\nu} \iff g_{\mu\nu}$$

$$S_{eff}(g, \psi_{SM}, \Lambda_X, \Lambda_U)$$

$$? \iff R_{\mu\nu}(g)$$

.....

“curvature” must be emergent!

- ▶ ? = what? Is it Λ_X ? or Λ_U ? or some other scale?

Gauge Invariance & Emergent Curvature

It proves efficacious to start with the trivial identity:

$$\delta S_V = \delta S_V + I_V - I_V$$

The diagram shows the decomposition of the variation of the action δS_V into two terms. Two curved arrows originate from the terms I_V and $-I_V$ in the equation above. The arrow from I_V points to the first integral on the right, and the arrow from $-I_V$ points to the second integral on the right.

$$\int d^4x \sqrt{-\eta} \frac{c_V}{2} \text{Tr}[V_{\mu\nu} V^{\mu\nu}]$$
$$\int d^4x \sqrt{-\eta} c_V (\Lambda_U^2 - \Lambda_X^2) \text{Tr}[V_\mu V^\mu]$$

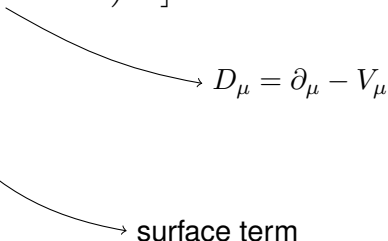
Gauge Invariance & Emergent Curvature

$$= \int d^4x \sqrt{-\eta} c_V (\Lambda_V^2 - \Lambda_X^2) \text{Tr}[V_\mu V^\mu]$$

$$+ \int d^4x \sqrt{-\eta} c_V \text{Tr}\left[V^\mu \left(-D^2 \eta_{\mu\nu} + D_\mu D_\nu + V_{\mu\nu}\right) V^\nu\right]$$

$$+ \int d^4x \sqrt{-\eta} c_V \text{Tr}\left[\partial_\nu (V_\mu V^{\mu\nu})\right]$$

$$- \int d^4x \sqrt{-\eta} \frac{c_V}{2} \text{Tr}\left[V_{\mu\nu} V^{\mu\nu}\right]$$



$D_\mu = \partial_\mu - V_\mu$

surface term

Gauge Invariance & Emergent Curvature

$$\frac{\eta_{\mu\nu} \rightleftharpoons g_{\mu\nu}}{\text{"Equivalence Principle"} \rightarrow}$$

$$\begin{aligned} & \int d^4x \sqrt{-g} c_V (\Lambda_U^2 - \Lambda_X^2) \text{Tr}[V_\mu V^\mu] \\ & + \int d^4x \sqrt{-g} c_V \text{Tr}[V^\mu (-\mathcal{D}^2 g_{\mu\nu} + \mathcal{D}_\mu \mathcal{D}_\nu + V_{\mu\nu}) V^\nu] \\ & + \int d^4x \sqrt{-g} c_V \text{Tr}[\nabla_\nu (V_\mu V^{\mu\nu})] \\ & - \int d^4x \sqrt{-g} \frac{c_V}{2} \text{Tr}[V_{\mu\nu} V^{\mu\nu}] \end{aligned}$$

$\mathcal{D}_\mu = \nabla_\mu - V_\mu$

surface term

Gauge Invariance & Emergent Curvature

$$= \int d^4x \sqrt{-g} c_V \text{Tr} \left[V^\mu \left(-\mathcal{D}^2 g_{\mu\nu} + \mathcal{D}_\mu \mathcal{D}_\nu + V_{\mu\nu} + \left(\Lambda_U^2 - \Lambda_X^2 \right) g_{\mu\nu} \right) V^\nu \right]$$

$$+ \int d^4x \sqrt{-g} c_V \text{Tr} \left[\nabla_\nu (V_\mu V^{\mu\nu}) \right]$$

$$- \int d^4x \sqrt{-g} \frac{c_V}{2} \text{Tr} \left[V_{\mu\nu} V^{\mu\nu} \right]$$



What if this is a fixed value assigned to **curvature**?



What if **curvature emerges** as $(\Lambda_U^2 - \Lambda_X^2) g_{\mu\nu} \iff R_{\mu\nu}(g)$?

Gauge Invariance & Emergent Curvature

$$\frac{(\eta_{\mu\nu} \mapsto g_{\mu\nu}) \& ((\Lambda_U^2 - \Lambda_X^2) g_{\mu\nu} \mapsto R_{\mu\nu}(g))}{\text{“Curved Equivalence Principle”}}$$

“Curved Equivalence Principle”

$$\begin{aligned} & \int d^4x \sqrt{-g} c_V \text{Tr} \left[V^\mu \left(-\mathcal{D}^2 g_{\mu\nu} + \mathcal{D}_\mu \mathcal{D}_\nu + V_{\mu\nu} + R_{\mu\nu}(g) \right) V^\nu \right] \\ & + \int d^4x \sqrt{-g} c_V \text{Tr} \left[\nabla_\nu (V_\mu V^{\mu\nu}) \right] \\ & - \int d^4x \sqrt{-g} \frac{c_V}{2} \text{Tr} \left[V_{\mu\nu} V^{\mu\nu} \right] \\ & = \int d^4x \sqrt{-g} \frac{c_V}{2} \text{Tr} \left[V_{\mu\nu} V^{\mu\nu} \right] - \int d^4x \sqrt{-g} \frac{c_V}{2} \text{Tr} \left[V_{\mu\nu} V^{\mu\nu} \right] \equiv 0 \quad \text{😊} \end{aligned}$$

if c_V is held unaffected

Gauge Invariance & Emergent Curvature

In summary:

$$\delta S_V = \delta S_V + I_V - I_V \xrightarrow[\substack{(\Lambda_U^2 - \Lambda_X^2)g_{\mu\nu} \iff R_{\mu\nu}(g)}]{\eta_{\mu\nu} \Rightarrow g_{\mu\nu}} 0$$

$\delta M_{W,Z}^2$ are **logarithmic**

EWB remains **spontaneous!**

photon and gluon
are both **massless**


Color & EM
are restored!

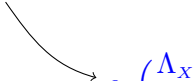
SM/UV Hierarchy Is Preserved

Holding $c_V\left(\frac{\Lambda_X}{\Lambda_U}\right)$ unaffected while $(\Lambda_{UV}^2 - \Lambda_{SM}^2)g_{\mu\nu} \iff R_{\mu\nu}(g)$



the hierarchy $\frac{\Lambda_X}{\Lambda_U}$ is preserved


$$c_O\left(\frac{\Lambda_X}{\Lambda_U}\right)$$


$$c_H\left(\frac{\Lambda_X}{\Lambda_U}\right)$$

Power-Law UV Contributions Set Curvature Sector

$$\delta S_{\text{power}} \xrightarrow[\substack{c_O, c_H \text{ held unchanged}}]{(\Lambda_U^2 - \Lambda_X^2) g_{\mu\nu} \mapsto R_{\mu\nu}(g)}$$

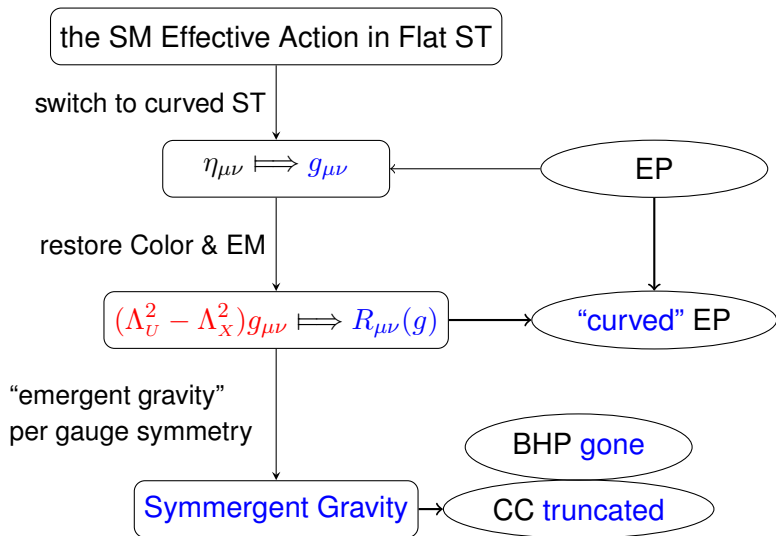
$$- \int d^4x \sqrt{-g} \underbrace{\frac{c_O}{4} (\Lambda_U^2 + \Lambda_X^2)}_{\frac{M_{Pl}^2}{2}} R(g)$$

quartic contribution
to vacuum energy
turns into **EH term!**

$$- \int d^4x \sqrt{-g} \frac{c_H}{4} H^\dagger H R(g)$$

big hierarchy
problem is **gone!**

“Symmergent Gravity”



Symmergent Gravity vs. Induced Gravity

Table 1: Contrasting symmergent gravity with Sakharov's induced gravity for $M_{Pl} \cong \Lambda_U$.

	Gravity Sector	Cosmological Constant	Higgs Mass	Color & EM
Induced Gravity	EH + HC	$\Lambda_U^4/M_{Pl}^2 \simeq M_{Pl}^2$	$\simeq M_{Pl}$	Broken
Symmergent Gravity	EH	$\Lambda_{EW}^4/M_{Pl}^2 \simeq m_\nu^2$	$\simeq \Lambda_{EW}$	Exact

Symmergent Gravity Results In Dim. Reg.

Holding $\log \frac{\Lambda_X}{\Lambda_U}$ unaffected while $(\Lambda_{UV}^2 - \Lambda_{SM}^2)g_{\mu\nu} \iff R_{\mu\nu}(g)$



$$\log \frac{\Lambda_X}{\Lambda_U} \rightsquigarrow -\frac{1}{\epsilon} - \log \frac{\mu}{\Lambda_X}$$

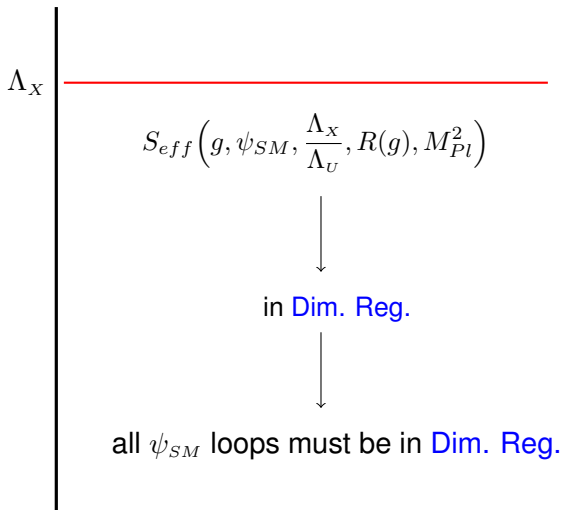


$$\delta S_{\log} \left(\psi_{SM}, \log \frac{\Lambda_X}{\Lambda_U} \right) \rightsquigarrow \text{Dim. Reg.}$$

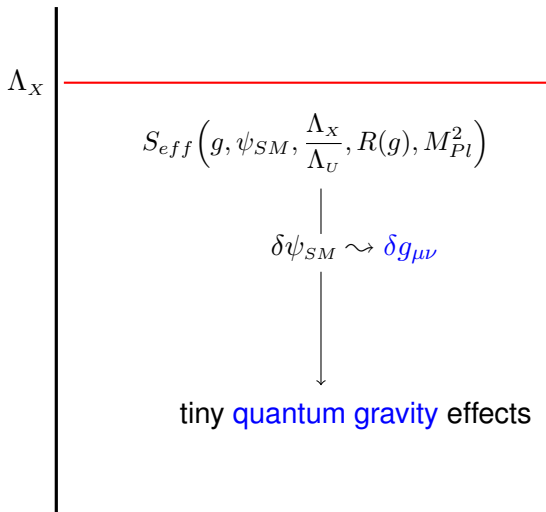


the usual RGEs

Symmergent Gravity Entails Dim. Reg. In The IR



Symmergent Gravity Is Only Feebly Quantum



Alas! SM Alone Can't Lead To Proper Gravity

In the SM (at one loop):

$$c_O \cong \frac{(n_b - n_f)}{64\pi^2} = \frac{-62}{64\pi^2}$$

not eligible for
inducing gravity

introduce **new fields**
to make $c_O \gtrsim 1$
so that $\Lambda_U \lesssim M_{Pl}$

A BSM Sector Is Required

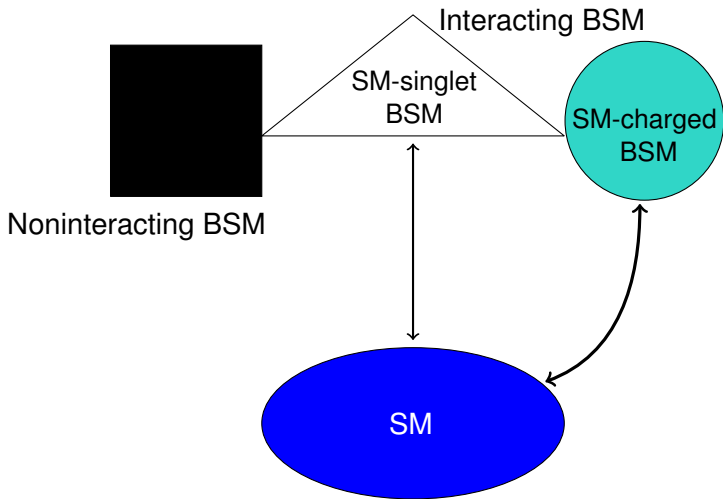
- ▶ All it has to do is to provide some

$$n_b^{BSM} - n_f^{BSM} \gtrsim 128\pi^2 + 62 \approx 1325$$

more bosons than fermions to insure $\Lambda_U \lesssim M_{Pl}$.

- ▶ Smaller the Λ_U/M_{Pl} more crowded the BSM.
- ▶ It **does not have to interact** with the SM but **it can**.

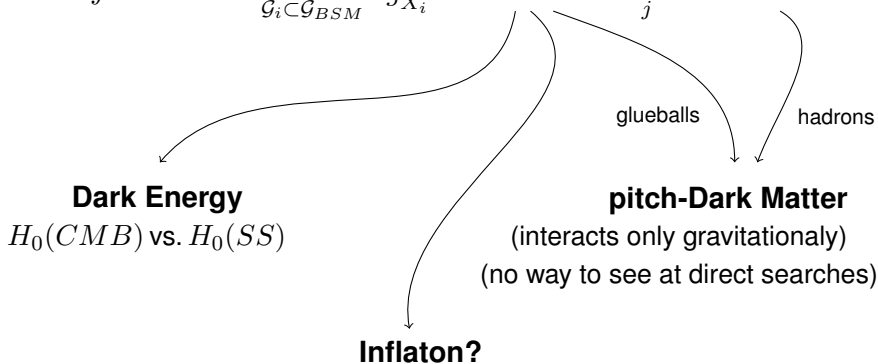
BSM Comes In Two Types



NonInt-BSM Is Sweet Home For “Dark Stuff”

- ▶ It can span only SM-singlet **non-Abelian gauge fields** and **fermions**.
- ▶ It enjoys gauge groups like $\mathcal{G}_{BSM} = SO(51), SU(26), E(8)^3, \dots$
- ▶ It is described by the action:

$$S^{(ni)} = \int d^4x \sqrt{-g} \left\{ - \sum_{\mathcal{G}_i \subset \mathcal{G}_{BSM}} \frac{1}{2g_{X_i}^2} \text{Tr} \{ X_{\mu\nu}^i X_i^{\mu\nu} \} + \sum_j \bar{\chi}_j (i\not{D} - m_{\chi_j}) \chi_j \right\}$$



SM-Charged Int-BSM Seems Nonexistent

- ▶ **No evidence** for interacting BSM at the LHC and DM searches.
- ▶ Interacting BSM fields with SM charges, if any, must weigh near Λ_{EW} since their loop contributions to Higgs boson mass

$$\delta m_H^2 \propto m_{BSM}^2 \log \frac{m_{BSM}}{\Lambda_U}$$

can be unacceptably large unless $m_{BSM} \cong \Lambda_{EW}$. But, this domain has already been **scanned** by the LHC and DM searches only to find **nothing**.

- ▶ It is for this reason that SUSY, Extra Dimensions, Compositeness and their various derivatives are **disfavored** by the data!

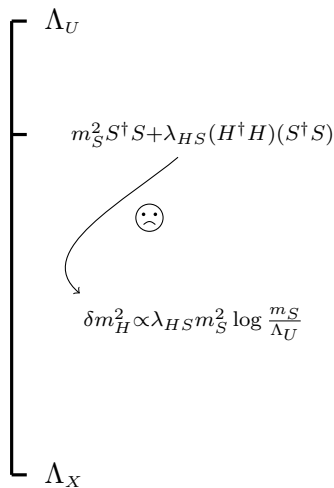
SM-Singlet Int-BSM Must Couple Seesawly

- ▶ **No evidence** for interacting BSM at the LHC and DM searches.
- ▶ But, it is **required** at least for Flavor, Strong CP, Baryogenesis, Inflation, whose models can contain scalars S .

- ▶ δm_H^2 can be suppressed only if

$$\lambda_{HS} \lesssim \Lambda_{EW}^2 / m_S^2$$

- ▶ $\delta \lambda_{HS} \propto \lambda_{HS} \implies \lambda_{HS}$ is natural

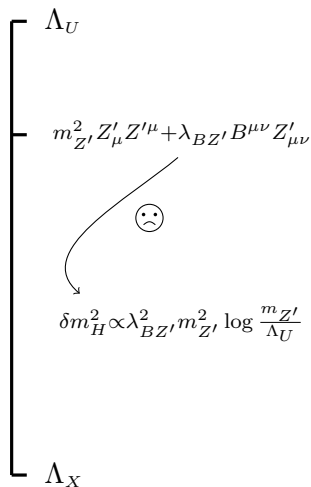


SM-Singlet Int-BSM Must Couple Seesawly

- ▶ Interacting BSM can contain a massive Z' gauge boson.
- ▶ It can mix with hypercharge gauge boson B kinetically.
- ▶ δm_H^2 can be suppressed only if


$$\lambda_{BZ'}^2 \lesssim \Lambda_{EW}^2 / m_{Z'}^2$$

- ▶ $\delta \lambda_{BZ'} \propto \lambda_{BZ'} \implies \lambda_{BZ'}$ is natural



SM-Singlet Int-BSM Must Couple Seesawly

- ▶ Neutrinos can be Dirac or Majorana.
- ▶ Dirac neutrinos **do not** destabilize the Higgs sector.
- ▶ Majorana neutrinos **do** destabilize with RH neutrinos N of mass m_N .
- ▶ A crowded noninteracting BSM can set $\Lambda_U = m_N$ and stabilize the Higgs.

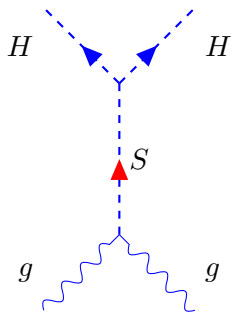
$$\Lambda_U$$
$$m_N N^T N + (\lambda_{LN} \bar{L} H N + h.c.)$$

$$\delta m_H^2 \propto \lambda_{LN}^2 m_N^2 \log \frac{m_N}{\Lambda_U}$$
$$\Lambda_X$$

It Is OK Even If LHC Discovers Nothing New

Symmergent gravity necessitates “no new interacting particles”:

- ▶ No new particle has to be **detected** at the LHC and others.
- ▶ No new particle has to be **detected** in DM searches.
- ▶ “SM+Noninteracting BSM” can **account for** all the data.

Heavier The BSM Higher The Luminosity Needed



Discovering a scalar S of width Γ_S and mass m_S requires a collider of center-of-mass energy $\sqrt{s} \simeq m_S$ and integrated luminosity $\simeq \left(\frac{m_S}{m_H}\right)^4$

$$\# \text{extra events} \simeq \frac{m_H^2}{m_S^2} \times \frac{m_H^2}{\Gamma_S^2} \times \# \text{SM events}$$

CCP Is “The” Problem

- ▶ The real **challenge** is to understand how CC can be reduced from Λ_{BSM}^4/M_{Pl}^2 down to H_0^2 .

- ▶ Solution of the CCP can **reveal** the BSM since

$$CC_{SM} + CC_{BSM} \cong H_0^2$$

can necessitate strong correlations with the SM.

- ▶ The CCP may well be a **door** to sought new physics.
(BSM, if noninteracting, can lie at ..., TeV's, GeV's, MeV's, ...)

We should ...

- ▶ investigate signals with no **theoretical prejudice**.
- ▶ keep in mind that we can end up with **no signal**.
- ▶ focus on **the CCP** to reveal structure of the BSM.
- ▶ explore **trans-UV** physics in regard to SM+BSM setup.

Thank You For Your Attention

References:

- ▶ *Naturalizing Gravity of the Quantum Fields, and the Hierarchy Problem*, [arXiv:1703.05733](https://arxiv.org/abs/1703.05733)
- ▶ *Curvature-Restored Gauge Invariance and Ultraviolet Naturalness*, [arXiv:1605.00377](https://arxiv.org/abs/1605.00377)
- ▶ *A Mechanism of Ultraviolet Naturalness*, [arXiv:1510.05570](https://arxiv.org/abs/1510.05570)

Acknowledgement

This work was supported through TUBITAK grant 115F212.